

## SPECTRAL EXPANSION FORMULA FOR A DISCONTINUOUS STURM-LIOUVILLE PROBLEM

KHANLAR R. MAMEDOV

*In memory of M. G. Gasymov on his 75th birthday*

**Abstract.** In this paper, a self adjoint boundary value problem with a piecewise continuous coefficient on the positive half line  $[0, \infty)$  is considered. The resolvent operator is constructed and the expansion formula with respect to eigenfunctions is obtained.

### 1. Introduction

On the semi axis  $0 < x < \infty$  we consider the boundary value problem by the differential equation

$$-y'' + q(x)y = \lambda^2 \rho(x)y \quad (1.1)$$

and the boundary condition

$$-[\alpha_1 y(0) - \alpha_2 y'(0)] = \lambda^2 [\beta_1 y(0) - \beta_2 y'(0)], \quad (1.2)$$

where  $\lambda$  is the spectral parameter and  $q(x)$  is the real function satisfying the condition

$$\int_0^\infty (1+x) |q(x)| dx < \infty \quad (1.3)$$

and

$$\rho(x) = \begin{cases} \alpha^2, & 0 \leq x < a, \\ 1, & x > a. \end{cases}$$

Here  $\alpha_i, \beta_i$  ( $i = 1, 2$ ) are real numbers and  $\gamma = \alpha_1 \beta_2 - \beta_1 \alpha_2 > 0$ .

In this paper, we obtain the expansion formula for the (1)-(3) boundary value problem in the half line. Similar problem for classical Sturm-Liouville equation, i.e. when  $\rho(x) = 1$  is studied in [15, 6, 13, 14, 2, 3]. When  $\rho(x) \neq 1$  direct and inverse problems of spectral analysis are investigated in Gasymovs' (see [4]) and his students works (see [1, 5, 7, 8, 9]).

In this study we used the integral expression for the Jost solution of the equation (1.1) as in [5] and differently from the classical boundary conditions, a boundary condition with dependence of eigenvalue has taken into account.

---

2010 *Mathematics Subject Classification.* 34L10, 34L40.

*Key words and phrases.* Sturm-Liouville operator, eigenvalue, eigenfunction, resolvent operator, expansion formula.

In this work, operator theoretic formulation is given for the boundary value problem (1.1), (1.2); resolvent operator is constructed and expansion formula with respect to eigenfunctions is obtained.

It is known from [5] that for all  $\lambda$  from the closed upper half plane equation (1.1) has a unique solution  $f(x, \lambda)$  which satisfies the condition

$$\lim_{x \rightarrow \infty} f(x, \lambda) e^{-i\lambda x} = 1$$

and that can be represented in the form

$$f(x, \lambda) = f_0(x, \lambda) + \int_{\mu^+(x)}^{+\infty} K(x, t) e^{i\lambda t} dt, \quad (1.4)$$

where

$$f_0(x, \lambda) = \begin{cases} e^{i\lambda x}, & x \geq a, \\ \frac{1}{2} \left( 1 + \frac{1}{\sqrt{\rho(x)}} \right) e^{i\lambda \mu^+(x)} + \frac{1}{2} \left( 1 - \frac{1}{\sqrt{\rho(x)}} \right) e^{i\lambda \mu^-(x)}, & 0 \leq x \leq a, \end{cases}$$

is the Jost solution of equation (1.1) when  $q(x) \equiv 0$  and

$$\mu^\pm(x) = \pm x \sqrt{\rho(x)} + a(1 \mp \sqrt{\rho(x)}).$$

Moreover, the kernel  $K(x, t)$  satisfies the estimate

$$\int_{\mu^+(x)}^{+\infty} |K(x, t)| dt \leq c \int_x^{+\infty} t |q(t)| dt, \quad 0 < c = \text{const.}$$

For real  $\lambda \neq 0$ , the functions  $f(x, \lambda)$ ,  $\overline{f(x, \lambda)}$  form the fundamental system of solutions of equation (1.1) and the Wronskian of this system is equal to  $2i\lambda$ :

$$W \left\{ f(x, \lambda), \overline{f(x, \lambda)} \right\} = f'(x, \lambda) \overline{f(x, \lambda)} - f(x, \lambda) \overline{f'(x, \lambda)} = 2i\lambda.$$

By  $w(x, \lambda)$  we denote the solutions of equation (1.1) satisfying the initial data

$$w(0, \lambda) = \alpha_2 + \lambda^2 \beta_2, \quad w'(0, \lambda) = \alpha_1 + \lambda^2 \beta_1$$

and define

$$\varphi(\lambda) \equiv [\alpha_1 f(0, \lambda) - \alpha_2 f'(0, \lambda)] - \lambda^2 [\beta_1 f(0, \lambda) - \beta_2 f'(0, \lambda)].$$

Using (1.4), it can be seen that  $\varphi(\lambda)$  may have only a finite number of zeros in the half plane  $\text{Im} \lambda > 0$ . Moreover, all these zeros are simple and lie on the imaginary axis. These assertions can be proved with the same method in Lemma 3.1.6 in [13] (see [9]).

## 2. The operator theoretic formulation of the boundary value problem (1.1), (1.2)

In the Hilbert space  $H_\rho \equiv L_2(0, \infty; \rho(x)) \times \mathbb{C}$ , let us define an inner product by

$$(F, G) := \int_0^\infty F_1(x) \overline{G_1(x)} \rho(x) dx + \frac{1}{\gamma} F_2 \overline{G_2}$$

for

$$F = \begin{pmatrix} F_1(x) \\ F_2 \end{pmatrix}, \quad G = \begin{pmatrix} G_1(x) \\ G_2 \end{pmatrix} \in H_\rho.$$

For convenience we put

$$F_2 = \beta_1 F_1(0) - \beta_2 F_1'(0) \equiv R'_0[F_1]$$

and define the operator  $L$  with

$$LF = \begin{pmatrix} l(F_1) \\ R_0(F_1) \end{pmatrix},$$

in the domain

$$D(L) = \{F \in H_\rho : F_1(x), F_1'(x) \in AC[0, b] \text{ for every } b, l(F_1) \in L_2(0, +\infty), \\ F_2 = \beta_1 F_1(0) - \beta_2 F_1'(0) \equiv R_0(F_1)\},$$

where

$$l(F_1) = \frac{1}{\rho(x)} \{-F_1'' + q(x)F_1\}.$$

We note that, the operator  $L$  is selfadjoint in the Hilbert space  $H_\rho$  and the boundary value problem (1.1), (1.2) is equivalent to the equation  $LF = \lambda^2 F$ .

### 3. Resolvent operator

If we assume that  $\lambda^2$  is not a spectrum point of the operator  $L$  we can easily say that the resolvent operator  $R_{\lambda^2}(L) = (L - \lambda^2 I)^{-1}$  exists. Now, let us find the expression of the operator  $R_{\lambda^2}(L)$ .

**Theorem 3.1.** *All numbers of the form  $\lambda^2$  ( $Im\lambda > 0$ ,  $\varphi(\lambda) \neq 0$ ) belong to the resolvent set of the operator  $L$ . The resolvent  $R_{\lambda^2}$  is the integral operator*

$$R_{\lambda^2}(L)F = \begin{pmatrix} (\tilde{G}_{x,x}, \bar{F})_{H_\rho} \\ R'_0 \left[ (\tilde{G}_{x,\lambda}, \bar{F})_{H_\rho} \right] \end{pmatrix}, \quad F \in H_\rho,$$

where

$$\tilde{G}_{x,\lambda} := \begin{pmatrix} G(x, \cdot, \lambda) \\ R'_0[G(x, \cdot, \lambda)] \end{pmatrix} = \begin{pmatrix} G(x, \cdot, \lambda) \\ -\frac{\gamma}{\varphi(\lambda)} f(x, \lambda) \end{pmatrix},$$

$$G(x, y, \lambda) = -\frac{1}{\varphi(\lambda)} \begin{cases} w(x, \lambda)f(y, \lambda), & x \leq y < +\infty, \\ f(x, \lambda)w(y, \lambda), & 0 \leq y < x. \end{cases}$$

*Proof.* Assume that  $F = \begin{pmatrix} F_1(x) \\ F_2 \end{pmatrix} \in D(L)$  and  $F_1(x)$  is zero in exterior of every interval. To construct the resolvent operator of  $L$  we need to solve the boundary problem

$$-y'' + q(x)y = \lambda^2 \rho(x)y + \rho(x)F_1(x), \quad (3.1)$$

$$-[\alpha_1 y(0) - \alpha_2 y'(0)] = \lambda^2 [\beta_1 y(0) - \beta_2 y'(0)] + F_2. \quad (3.2)$$

Let us find the solutions of the problem (3.1), (3.2) which have the form

$$y(x, \lambda) = c_1(x, \lambda)w(x, \lambda) + c_2(x, \lambda)f(x, \lambda) \quad (3.3)$$

by noting that the functions  $w(x, \lambda)$  and  $f(x, \lambda)$  are the solutions of homogeneous problem for  $Im\lambda > 0$ .

By applying the method of variation of constants, we have the system of equations

$$\begin{cases} c_1'(x, \lambda)w(x, \lambda) + c_2'(x, \lambda)f(x, \lambda) = 0, \\ c_1'(x, \lambda)w'(x, \lambda) + c_2'(x, \lambda)f'(x, \lambda) = -\rho(x)F_1(x). \end{cases} \quad (3.4)$$

Since  $y(x, \lambda) \in L_{2,\rho}(0, \infty)$ , then  $c_1(\infty, \lambda) = 0$ . Using this relation and the system of equations (3.4) we get

$$\begin{aligned} c_1(x, \lambda) &= - \int_x^\infty \frac{f(t, \lambda)}{\varphi(\lambda)} F_1(t) \rho(t) dt, \\ c_2(x, \lambda) &= c_2(0, \lambda) - \int_0^x \frac{w(t, \lambda)}{\varphi(\lambda)} F_1(t) \rho(t) dt. \end{aligned} \quad (3.5)$$

Taking into consideration (3.5) into (3.3) we obtain

$$y(x, \lambda) = \int_0^\infty G(x, t, \lambda) F_1(t) \rho(t) dt + c_2(0, \lambda) f(x, \lambda),$$

where

$$G(x, t, \lambda) = \frac{1}{\varphi(\lambda)} \begin{cases} w(x, \lambda) f(y, \lambda), & x \leq y < +\infty, \\ f(x, \lambda) w(y, \lambda), & 0 \leq y < x. \end{cases}$$

Using formula (3.2) we have  $c_2(0, \lambda) = \frac{F_2}{\varphi(\lambda)}$ , thus

$$y(x, \lambda) = \int_0^\infty G(x, t, \lambda) F_1(t) \rho(t) dt + \frac{F_2}{\varphi(\lambda)} f(x, \lambda) = \left( \tilde{G}_{x,\lambda}, \bar{F} \right)_{H_\rho},$$

where

$$\tilde{G}_{x,\lambda} = \begin{pmatrix} G(x, \cdot, \lambda) \\ R'_0 [G(x, \cdot, \lambda)] \end{pmatrix}.$$

Theorem is proved.  $\square$

**Lemma 3.1.** *Let the function  $F_1(x)$  is continuously differentiable two times and finite at infinity. Then as  $|\lambda| \rightarrow \infty$ ,  $\text{Im} \lambda > 0$  the following equation holds:*

$$\left( \tilde{G}, \tilde{F} \right) = \int_0^\infty G(x, t, \lambda) F_1(t) \rho(t) dt + \frac{F_2}{\varphi(\lambda)} f(x, \lambda) = -\frac{F_1(x)}{\lambda^2} + O\left(\frac{1}{\lambda^2}\right). \quad (3.6)$$

*Proof.* Using Theorem 1 and integrating by parts, we write

$$\begin{aligned} \int_0^\infty G(x, t, \lambda) F_1(t) \rho(t) dt &= - \int_x^\infty \frac{w(x, \lambda) f(t, \lambda)}{\varphi(\lambda)} F_1(t) \rho(t) dt - \\ &- \int_x^\infty \frac{w(t, \lambda) f(x, \lambda)}{\varphi(\lambda)} F_1(t) \rho(t) dt = \\ &= \frac{1}{\lambda^2} \int_x^\infty \frac{w(x, \lambda)}{\varphi(\lambda)} \{f''(t, \lambda) - q(t) f(t, \lambda)\} F_1(t) dt + \\ &+ \frac{1}{\lambda^2} \int_x^\infty \frac{f(x, \lambda)}{\varphi(\lambda)} \{w''(t, \lambda) - q(t) w(t, \lambda)\} F_1(t) dt = \\ &= \frac{W\{w(x, \lambda), f(x, \lambda)\}}{\lambda^2 \varphi(\lambda)} F_1(x) + \\ &+ \frac{1}{\lambda^2} [w(0, \lambda) F_1'(0) - w'(0, \lambda) F_1(0)] - \\ &- \frac{1}{\lambda^2} \int_0^\infty G(x, t, \lambda) \tilde{F}_1(t) dt, \end{aligned}$$

where

$$\tilde{F}_1(t) = F_1''(t) + q(t)F_1(t).$$

Thus (3.6) holds. Lemma is proved.  $\square$

#### 4. Expansion formula

Put

$$F(x, \lambda) = \begin{cases} \left( \tilde{G}_{x, \lambda}, \bar{F} \right), & \text{Im} \lambda \geq 0, \\ \left( \tilde{G}_{x, \bar{\lambda}}, \bar{F} \right), & \text{Im} \lambda < 0. \end{cases}$$

Let us  $\Gamma_R$  denote the circle of radius  $R$  and center is zero which boundary contour is positive oriented. Let us  $\Gamma_{R, \epsilon}$  denote boundary contour positive oriented in plane  $D = \{z : |z| \leq R, \text{Im} z \geq \epsilon\}$  and  $\Gamma_{R', \epsilon}$  denote boundary contour negative oriented in the plane  $D = \{z : |z| \leq R, \text{Im} z < \epsilon\}$ . Then we can use the property of the integration

$$\int_{\Gamma_{R, \epsilon}} = \int_{\Gamma_R} + \int_{\Gamma_{R', \epsilon}}. \quad (4.1)$$

Now multiplying both sides of equality (3.6) by  $\frac{1}{2\pi i} \lambda$  and integrating over  $\lambda$  the contour  $\Gamma_{R, \epsilon}$  we obtain

$$\frac{1}{2\pi i} \int_{\Gamma_{R, \epsilon}} \lambda F(x, \lambda) d\lambda = -\frac{1}{2\pi i} \int_{\Gamma_R} \frac{F_1(x)}{\lambda} d\lambda + \frac{1}{2\pi i} \int_{\Gamma_{R', \epsilon}} O\left(\frac{1}{\lambda}\right) d\lambda. \quad (4.2)$$

According to the equation (4.1) we get

$$\frac{1}{2\pi i} \int_{\Gamma_{R, \epsilon}} \lambda F(x, \lambda) d\lambda = \frac{1}{2\pi i} \int_{\Gamma_R} \lambda F(x, \lambda) d\lambda + \frac{1}{2\pi i} \int_{\Gamma_{R', \epsilon}} \lambda F(x, \lambda) d\lambda. \quad (4.3)$$

Using (4.2) let us calculate the integral on the right hand side

$$\frac{1}{2\pi i} \int_{\Gamma_R} \lambda F(x, \lambda) d\lambda = -\frac{1}{2\pi i} \int_{\Gamma_R} \frac{F_1(x)}{\lambda} d\lambda + \frac{1}{2\pi i} \int_{\Gamma_R} O\left(\frac{1}{\lambda}\right) d\lambda \rightarrow -F_1(x) \quad (R \rightarrow \infty),$$

taking into account (4.3) we have

$$\begin{aligned} \lim_{\epsilon \rightarrow 0, R \rightarrow \infty} \frac{1}{2\pi i} \int_{\Gamma_{R, \epsilon}} \lambda F(x, \lambda) d\lambda &= -F_1(x) \\ &+ \frac{1}{2\pi i} \int_{-\infty}^{+\infty} \lambda [F(x, \lambda + i0) - F(x, \lambda - i0)] d\lambda. \end{aligned}$$

On the other hand using the residue calculus we get

$$\frac{1}{2\pi i} \int_{\Gamma_{R, \epsilon}} \lambda F(x, \lambda) d\lambda = \sum_{j=1}^n \text{Res}_{\lambda=i\lambda_j} [\lambda F(x, \lambda)] + \sum_{j=1}^n \text{Res}_{\lambda=-i\lambda_j} [\lambda F(x, \lambda)].$$

From the last two relations we get

$$\begin{aligned} F_1(x) = & - \sum_{j=1}^n \text{Res}_{\lambda=i\lambda_j} [\lambda F(x, \lambda)] + \sum_{j=1}^n \text{Res}_{\lambda=-i\lambda_j} [\lambda F(x, \lambda)] + \\ & + \frac{1}{2\pi i} \int_{-\infty}^{+\infty} \lambda [F(x, \lambda + i0) - F(x, \lambda - i0)] d\lambda. \end{aligned} \quad (4.4)$$

Let  $\psi(x, \lambda)$  be the solution of (1.1) satisfying the initial conditions

$$\psi(0, \lambda) = \beta_2, \quad \psi'(0, \lambda) = \beta_1.$$

It is clear that  $W\{w(x, \lambda), \psi(x, \lambda)\} = \gamma > 0$ . Then for  $Im\lambda > 0$ ,  $f(x, \lambda)$  is written as the linear combination of the solutions  $w(x, \lambda)$  and  $\psi(x, \lambda)$ , i.e.

$$f(x, \lambda) = c_1 w(x, \lambda) + c_2 \psi(x, \lambda).$$

Since  $W\{w(x, \lambda), \psi(x, \lambda)\} = c_1 \gamma$ ,  $W\{w(x, \lambda), f(x, \lambda)\} = c_2 \gamma$

$$f(x, \lambda) = \frac{\beta_2 f'(0, \lambda) - \beta_1 f(0, \lambda)}{\gamma} w(x, \lambda) + \frac{\varphi(\lambda)}{\gamma} \psi(x, \lambda).$$

Thus

$$\begin{aligned} G(x, t, \lambda) &= -\frac{\beta_2 f'(0, \lambda) - \beta_1 f(0, \lambda)}{\gamma \varphi(\lambda)} w(x, \lambda) w(t, \lambda) - \\ &\quad - \frac{1}{\gamma} \begin{cases} w(x, \lambda) \psi(t, \lambda), & x \leq t < +\infty, \\ \psi(x, \lambda) w(t, \lambda), & 0 \leq t < x. \end{cases} \end{aligned}$$

Then for  $Im\lambda > 0$  we obtain

$$\begin{aligned} F(x, \lambda) &= -\frac{\beta_2 f'(0, \lambda) - \beta_1 f(0, \lambda)}{\gamma \varphi(\lambda)} w(x, \lambda) \int_0^\infty w(t, \lambda) F_1(t) \rho(t) dt + \\ &\quad + \frac{\psi(x, \lambda)}{\gamma} \int_0^x w(t, \lambda) F_1(t) \rho(t) dt + \frac{1}{\gamma} \psi(x, \lambda) \int_0^x \psi(t, \lambda) F_1(t) \rho(t) dt + \\ &\quad + \frac{F_2}{\gamma \varphi(\lambda)} [\beta_2 f'(0, \lambda) - \beta_1 f(0, \lambda)] w(x, \lambda) - \frac{F_2}{\gamma} \psi(x, \lambda). \end{aligned} \quad (4.5)$$

It follows that

$$\begin{aligned} Res_{\lambda=i\lambda_j} [\lambda F(x, \lambda)] &= -\frac{i\lambda_j}{\varphi'(i\lambda_j)} f(0, i\lambda_j) w(x, i\lambda_j) \int_0^\infty w(t, i\lambda_j) F_1(t) \rho(t) dt + \\ &\quad + \frac{F_2 i\lambda_j}{\gamma \varphi'(i\lambda_j)} [\beta_2 f'(0, i\lambda_j) - \beta_1 f(0, i\lambda_j)] w(x, i\lambda_j). \end{aligned}$$

With the help of equalities (1.5) and (1.8) we get

$$\begin{aligned} \sum_{j=1}^n Res_{\lambda=i\lambda_j} [\lambda F(x, \lambda)] &+ \sum_{j=1}^n Res_{\lambda=-i\lambda_j} [\lambda F(x, \lambda)] = \\ &- m_j^2 f(x, i\lambda_j) \int_0^\infty f(t, i\lambda_j) F_1(t) \rho(t) dt + \\ &+ \frac{1}{\gamma} [\beta_2 f'(0, i\lambda_j) - \beta_1 f(0, i\lambda_j)] m_j^2 f(x, i\lambda_j) F_2 = \\ &= -(F, U_j(\lambda))_{H_\rho} m_j f(x, i\lambda_j), \end{aligned} \quad (4.6)$$

where

$$F = \begin{pmatrix} F_1(x) \\ F_2 \end{pmatrix}, \quad U_j(x) = m_j \begin{pmatrix} f(x, i\lambda_j) \\ \beta_2 f'(0, i\lambda_j) - \beta_1 f(0, i\lambda_j) \end{pmatrix}, \quad j = 1, \dots, n. \quad (4.7)$$

Now let us calculate

$$\frac{1}{2\pi i} \int_{-\infty}^{+\infty} \lambda [F(x, \lambda + i0) - F(x, \lambda - i0)] d\lambda.$$

From formula (3.5) and equality  $F(x, \lambda - i0) = \overline{F(x, \lambda + i0)}$  we have

$$\begin{aligned}
& F(x, \lambda + i0) - F(x, \lambda - i0) = \\
& \left[ \frac{\overline{\beta_2 f'(0, \lambda) - \beta_1 f(0, \lambda)}}{\gamma \overline{\varphi(\lambda)}} - \frac{\beta_2 f'(0, \lambda) - \beta_1 f(0, \lambda)}{\gamma \varphi(\lambda)} \right] w(x, \lambda) \int_0^\infty w(t, \lambda) F_1(t) \rho(t) dt + \\
& + \frac{F_2}{\gamma} \left[ \frac{\beta_2 f'(0, \lambda) - \beta_1 f(0, \lambda)}{\gamma \varphi(\lambda)} - \frac{\overline{\beta_2 f'(0, \lambda) - \beta_1 f(0, \lambda)}}{\gamma \overline{\varphi(\lambda)}} \right] w(x, \lambda) = \\
& = \frac{1}{\gamma} \left[ \frac{(\alpha_1 \beta_2 - \beta_1 \alpha_2) \left[ f'(0, \lambda) \overline{f(0, \lambda)} - f(0, \lambda) \overline{f'(0, \lambda)} \right]}{|\varphi(\lambda)|^2} \right] w(x, \lambda) \times \\
& \quad \times \int_0^\infty w(t, \lambda) F_1(t) \rho(t) dt + \\
& \quad + \frac{F_2 (\alpha_1 \beta_2 - \beta_1 \alpha_2)}{\gamma |\varphi(\lambda)|^2} \left[ f(0, \lambda) \overline{f'(0, \lambda)} - f'(0, \lambda) \overline{f(0, \lambda)} \right] w(x, \lambda) = \\
& = \frac{2i\lambda}{|\varphi(\lambda)|^2} w(x, \lambda) \int_0^\infty w(t, \lambda) F_1(t) \rho(t) dt - \frac{2i\lambda F_2}{|\varphi(\lambda)|^2} w(x, \lambda).
\end{aligned}$$

It follows that

$$\begin{aligned}
& \frac{1}{2\pi i} \int_{-\infty}^{+\infty} \lambda [F(x, \lambda + i0) - F(x, \lambda - i0)] d\lambda = \\
& = \frac{2}{\pi} \int_0^\infty \frac{\lambda^2}{|\varphi(\lambda)|^2} w(x, \lambda) \int_0^\infty w(t, \lambda) F_1(t) \rho(t) dt d\lambda - \\
& - \frac{2F_2}{\pi} \int_0^\infty \frac{\lambda}{|\varphi(\lambda)|^2} w(x, \lambda) d\lambda = \int_0^\infty T(F, \lambda) \frac{1}{\sqrt{2\pi}} \left[ \overline{f(x, \lambda)} - S(\lambda) f(x, \lambda) \right] d\lambda,
\end{aligned} \tag{4.8}$$

where

$$T(F, \lambda) \equiv -\frac{1}{\sqrt{2\pi}} \int_0^\infty \frac{2i\lambda}{|\varphi(\lambda)|^2} w(t, \lambda) F_1(t) \rho(t) dt + F_2 \frac{2i\lambda \gamma^2}{\varphi(\lambda)}.$$

Taking (4.6) and (4.8) into (4.4) we get the expansion formula with respect to eigenfunctions as

$$F \equiv \begin{pmatrix} F_1(x) \\ F_2 \end{pmatrix} = \sum_{j=1}^n (F, U_j(x))_{H_\rho} U_j(x) + \int_0^\infty T(F, \lambda) U(x, \lambda) d\lambda,$$

where

$$U(x, \lambda) \equiv \frac{1}{\sqrt{2\pi}} \begin{pmatrix} \overline{f(x, \lambda)} - S(\lambda) f(x, \lambda) \\ -\frac{2i\lambda \gamma}{\varphi(\lambda)} \end{pmatrix} = \frac{1}{\sqrt{2\pi}} \begin{pmatrix} \frac{2i\lambda f(x, \lambda)}{\varphi(\lambda)} \\ -\frac{2i\lambda \gamma}{\varphi(\lambda)} \end{pmatrix}. \tag{4.9}$$

## References

- [1] A. A. Darwish, The inverse problem for a singular boundary value problem, *New Zealand Journal of Mathematics*, **22** (1993), 37–66.
- [2] G. Freiling, V. Yurko, *Inverse Sturm-Liouville problems*, NovaScience Publishers INC. 2008.
- [3] C. T. Fulton, Singular eigenvalue problems with eigenvalue parameter contained in the boundary conditions, *Proc. Ro. Soc. Edin.*, **87 A** (1980), 1-34.
- [4] M. G. Gasymov, The direct and inverse problem of spectral analysis for a class of equations with a discontinuous coefficient, in: *Non-Classical Methods in Geophysics*, M. M. Lavrent'ev, Ed., Nauka, Novosibirsk, Russia, (1977) 37–44.
- [5] I. M. Guseinov, R.T. Pashaev, On an inverse problem for a second order differential equation, *Usp. Math. Nauk*, **57(3)** (2002), 597–598.
- [6] B. M. Levitan, I. S. Sargsjan, *Introduction to spectral theory*, American Mathematical Society, 1975.
- [7] Kh. R. Mamedov, Uniqueness of the solution to the inverse problem of scattering theory for Sturm-Liouville operator with discontinuous coefficient, *Math. Notes*, **74(1)** (2003), 136–140.
- [8] Kh. R. Mamedov, Uniqueness of the solution of the inverse problem of scattering theory for Sturm-Liouville operator with discontinuous coefficient, *Proceedings of IMM of NAS Azerbaijan* **24** (2006), 263–272.
- [9] Kh. R. Mamedov, On an inverse scattering problem for a discontinuous Sturm-Liouville equation with a spectral parameter in the boundary condition, *Boundary Value Problems*, **2010**, 17 pages, doi 10.1155/2010/171967.
- [10] Kh. R. Mamedov, V. Ala, On the solution of a boundary value problem related to the heat transmission, *American Journal of Applied Mathematics*, **2(2)** (2014), 54–59, doi: 10.11648/j.ajam.20140202.12.
- [11] Kh. R. Mamedov, F. A. Cetinkaya, Boundary value problem for a Sturm-Liouville operator with piecewise continuous coefficient, *Hacettepe Journal of Mathematics and Statistics* (accepted)
- [12] Kh. R. Mamedov, N. Palamut, On a direct problem of scattering theory for a class of Sturm Liouville operator with discontinuous coefficient, *Proceedings of the Jangjeon Mathematical Society*, **12(2)** (2009), 243–251.
- [13] V. A. Marchenko, *Sturm-Liouville operators and their applications*, AMS Chelsea Publishing, 2011.
- [14] M. A. Naimark, *Linear differential equations*, Frederick Ungar Publishing, 1967.
- [15] E. C. Titchmarsh, *Eigenfunction expansions*, Oxford, 1962.

Khanlar R. Mamedov

*Department of Mathematics, Mersin University, Mersin, 33343, Turkey.*

*Institute of Mathematics and Mechanics, National Academy of Sciences of Azerbaijan, Baku, AZ1141, Azerbaijan.*

E-mail address: hanlar@mersin.edu.tr

Received: June 12, 2014; Accepted: July 14, 2014